The hemispherical power asymmetry after Planck

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The hemispherical power asymmetry



ESA and the Planck Collaboration 2013

The Dipole Modulation of CMB

$$\Delta T(\hat{n}) = \Delta T_{\rm iso}(\hat{n}) \left[1 + A\hat{\lambda} \cdot \hat{n} \right]$$

gives:

$$a_{T,\ell m} = \tilde{a}_{T,\ell m} + A \sum_{\ell'm'} \tilde{a}_{T,\ell'm'} \sqrt{(2\ell+1)(2\ell'+1)}$$

$$\times \left[\frac{\sqrt{(\ell-m)(\ell+m)}}{(2\ell+1)(2\ell-1)} \delta_{\ell'\ell-1} \delta_{m'm} + \frac{\sqrt{(\ell-m+1)(\ell+m+1)}}{(2\ell+1)(2\ell+3)} \delta_{\ell'\ell+1} \delta_{m'm} \right]$$

$$= \tilde{a}_{T,\ell m} + A\xi_{-}(\ell,m) \tilde{a}_{T,\ell-1m} + A\xi_{+}(\ell,m) \tilde{a}_{T,\ell+1m}$$

So:

$$\begin{aligned} \langle a_{T,\ell m} a_{T,\ell' m'}^* \rangle &= \tilde{C}_{\ell}^{TT} \delta_{\ell' \ell} \delta_{m'm} + [\tilde{C}_{\ell'}^{TT} + \tilde{C}_{\ell}^{TT}] \\ &\times \{ A\xi_{-}(\ell,m) \delta_{\ell' \ell-1} \delta_{m'm} + A\xi_{+}(\ell,m) \delta_{\ell' \ell+1} \delta_{m'm} \} \end{aligned}$$

The S_{H}^{TT} estimator

Define:

$$S_{H}^{TT} = \sum_{\ell=\ell_{\min}}^{\ell_{\max}} \frac{\ell(\ell+1)}{2\ell+1} \sum_{m} a_{T,\ell m} a_{T,\ell+1m}^{*}$$

Мар	S_H^{TT} in 10 ⁻² mK ²	Α	(<i>l</i> , <i>b</i>)	P-value
Comm	2.55 ± 0.68	0.082 ± 0.018	$(232^{\circ} \pm 18^{\circ}, -14^{\circ} \pm 18^{\circ})$	0.20%
SMICA(i)	2.39 ± 0.70	0.069 ± 0.013	$(236^{\circ} \pm 27^{\circ}, -11^{\circ} \pm 20^{\circ})$	0.70%
SMICA(f)	2.44 ± 0.71	0.078 ± 0.019	$(242^{\circ} \pm 16^{\circ}, -17^{\circ} \pm 20^{\circ})$	0.50%

with $\ell_{min} = 2$ and $\ell_{max} = 64$.

Scale Dependence

$$T_{TT} = \frac{S_{H}^{TT}}{\sum_{\ell=\ell_{min}}^{\ell_{max}} \ell(\ell+1)C_{\ell}^{TT}}$$

/ Range	S_H^T in 10 ⁻² mK ²	А	(I, b)	P-value	r _{TT}
2-64	2.55 ± 0.68	0.082 ± 0.018	(232°, -14°)	0.20%	0.065
30-64	1.00 ± 0.43	0.052 ± 0.019	(194°, -4°)	66.6%	0.040
30-100	0.91 ± 0.72	0.018 ± 0.011	(277°,4°)	83.0%	0.013



Ghosh, S. et al. 2016

Scale Dependence



CMB Polarization



The polarized CMB sky



The Ps, Es and Bs

Define:

$$P(\hat{n}) = Q(\hat{n}) + iU(\hat{n})$$
$$P^*(\hat{n}) = Q(\hat{n}) - iU(\hat{n})$$

On rotation: $(Q \pm iU)'(\hat{n}) = (Q \pm iU)(\hat{n})e^{\pm i2\psi}$

$$P(\hat{n}) = \sum_{\ell,m} a_{2,\ell m 2} Y_{\ell m}(\hat{n}) = -\sum_{\ell,m} (a_{E,\ell m} + ia_{B,\ell m})_2 Y_{\ell m}(\hat{n})$$

$$P^*(\hat{n}) = \sum_{\ell,m} a_{-2,\ell m - 2} Y_{\ell m}(\hat{n}) = -\sum_{\ell,m} (a_{E,\ell m} - ia_{B,\ell m})_{-2} Y_{\ell m}(\hat{n})$$

$$= -\sum_{\ell,m} (a_{E,\ell m}^* - ia_{B,\ell m}^*)_2 Y_{\ell m}^*(\hat{n}),$$

Modulation of CMB Polarization

$$(Q \pm iU)(\hat{n}) \propto \int dr \frac{d\tau}{dr} e^{-\tau(r)} \sum_{m} a_{T,2m}(r\hat{n})_{\pm 2} Y_{2m}(\hat{n})$$

Hu, W. 2000; Contreras, D. 2017

Since
$$a_{T,2m}(r_{ls}\hat{n}) = \tilde{a}_{T,2m}(r_{ls}\hat{n}) \left[1 + A\hat{\lambda} \cdot \hat{n} \right]$$
:
 $P(\hat{n}) = \tilde{P}(\hat{n}) \left(1 + A\hat{\lambda} \cdot \hat{n} \right).$

So:

$$a_{\pm 2,\ell m} = \tilde{a}_{\pm 2,\ell m} + A \sqrt{\frac{4\pi}{3}} \sum_{\ell' m'} \tilde{a}_{\pm 2,\ell' m'} \int_{\pm 2} Y_{\ell' m'}(\hat{n}) Y_{10}(\hat{n})_{\pm 2} Y_{\ell m}^{*}(\hat{n}) d\Omega.$$

The E/B-mode modulation

$$a_{E,\ell m} = \tilde{a}_{E,\ell m} + A\alpha_{-}\tilde{a}_{E,\ell-1m} + iA\alpha_{0}\tilde{a}_{B,\ell m} + A\alpha_{+}\tilde{a}_{E,\ell+1m}$$
$$a_{B,\ell m} = \tilde{a}_{B,\ell m} + A\alpha_{-}\tilde{a}_{B,\ell-1m} - iA\alpha_{0}\tilde{a}_{E,\ell m} + A\alpha_{+}\tilde{a}_{B,\ell+1m},$$

where,

$$\alpha_{-} = \frac{1}{\ell} \sqrt{\frac{(\ell-2)(\ell+2)(\ell-m)(\ell+m)}{(2\ell-1)(2\ell+1)}}$$

$$\alpha_{0} = \frac{2m}{\ell(\ell+1)}$$

$$\alpha_{+} = \frac{1}{\ell+1} \sqrt{\frac{(\ell-1)(\ell+3)(\ell-m+1)(\ell+m+1)}{(2\ell+1)(2\ell+3)}}.$$

The Correlations

The different auto correlations, written up to first order in A, are:

$$\begin{split} \langle a_{E,\ell m} a_{E,\ell'm'}^* \rangle &= \tilde{C}_{\ell}^{EE} \delta_{\ell'\ell} \delta_{m'm} + [\tilde{C}_{\ell'}^{EE} + \tilde{C}_{\ell}^{EE}] \\ &\times \{A\alpha_{-}(\ell,m)\delta_{\ell'\ell-1}\delta_{m'm} + A\alpha_{+}(\ell,m)\delta_{\ell'\ell+1}\delta_{m'm}\} \\ \langle a_{B,\ell m} a_{B,\ell'm'}^* \rangle &= \tilde{C}_{\ell}^{BB} \delta_{\ell'\ell} \delta_{m'm} + [\tilde{C}_{\ell'}^{BB} + \tilde{C}_{\ell}^{BB}] \\ &\times \{A\alpha_{-}(\ell,m)\delta_{\ell'\ell-1}\delta_{m'm} + A\alpha_{+}(\ell,m)\delta_{\ell'\ell+1}\delta_{m'm}\} \,. \end{split}$$

The P^2 map

With noise $N_P(\hat{n})$ and mask $W(\hat{n})$:

$$P_{obs} = \tilde{P}_{s}(\hat{n})W(\hat{n})\left(1 + A\hat{\lambda}\cdot\hat{n}\right) + N_{P}(\hat{n})W(\hat{n})$$

So:

$$\begin{split} |P_{\text{obs}}(\hat{n})|^2 &= |\tilde{P}_s(\hat{n})|^2 W^2(\hat{n}) \left(1 + 2A\hat{\lambda} \cdot \hat{n}\right) + \tilde{P}_s(\hat{n}) N_P^*(\hat{n}) W^2(\hat{n}) \left(1 + A\hat{\lambda} \cdot \hat{n}\right) \\ &+ \tilde{P}_s^*(\hat{n}) N_P(\hat{n}) W^2(\hat{n}) \left(1 + 2A\hat{\lambda} \cdot \hat{n}\right) + |N_P(\hat{n})|^2 W^2(\hat{n}) \end{split}$$

Ensemble average:

$$\langle |P_{\text{obs}}(\hat{n}_i)|^2 \rangle = W^2(\hat{n}_i) \sum_{\ell} \left(\frac{2\ell+1}{4\pi} \right) \left\{ \left[\bar{C}_{\ell}^{EE} + \bar{C}_{\ell}^{BB} \right] \left[1 + 2A\hat{\lambda} \cdot \hat{n}_i \right] \right. \\ \left. + \left[\bar{N}_{\ell}^{EE} + \bar{N}_{\ell}^{BB} \right] \right\}$$

The Harmonic Space Estimator

The S_{H}^{EE} statistic:

$$S_{H}^{EE} = \sum_{\ell=\ell_{\min}}^{\ell_{\max}} \frac{\ell(\ell+1)}{2\ell+1} \sum_{m} a_{E,\ell m} a_{E,\ell+1m}^{*}$$

The *r_{EE}* measure:

$$r_{EE} = \frac{S_{H}^{EE}}{\sum_{\ell=\ell_{min}}^{\ell_{max}} \ell(\ell+1)C_{\ell}^{EE}}$$

Weighted Pixel-space Direction Estimators

The weight factor:

$$\omega_j = |\cos \theta_j|$$

Weighted average:

$$\langle |P_{\text{obs}}(\hat{n})|^2 \rangle_w = \frac{\sum_j \omega_j |P_{\text{obs}}(\hat{n}_j)|^2}{\sum_{j'} W(\hat{n}_{j'}) \omega_{j'}}$$

The direction estimators:

$$\begin{aligned} R_{i}^{w} &= \frac{\langle |P_{obs}(\hat{n})|^{2} \rangle_{U_{i},w}}{\langle |P_{obs}(\hat{n})|^{2} \rangle_{L_{i},w}} \\ D_{i}^{w} &= \frac{\langle |P_{obs}(\hat{n})|^{2} \rangle_{U_{i},w} - \langle |P_{obs}(\hat{n})|^{2} \rangle_{L_{i},w}}{\langle |P_{obs}(\hat{n})|^{2} \rangle_{U_{i},w} + \langle |P_{obs}(\hat{n})|^{2} \rangle_{L_{i},w}} \end{aligned}$$

The pixel space amplitude estimator

$$\hat{A} = \frac{\left[\langle |P_{\text{obs}}(\hat{n})|^2 \rangle_U - \langle |P_{\text{obs}}(\hat{n})|^2 \rangle_L\right]_{\text{max}}}{K \sum_{\ell} \left(\frac{2\ell+1}{2\pi}\right) \left(\bar{C}_{\ell}^{EE} + \bar{C}_{\ell}^{BB}\right)}$$

$$\mathcal{K} = \left\{ \int_{U} W^{2}(\hat{n}) d\Omega \right\}^{-1} \int_{U} W^{2}(\hat{n}) \cos \theta d\Omega - \left\{ \int_{L} W^{2}(\hat{n}) d\Omega \right\}^{-1} \int_{L} W^{2}(\hat{n}) \cos \theta d\Omega.$$

Planck Polarization Data

Planck polarization data has significant contamination from systematics on large angular scales.

Source of contamination are the LFI calibration uncertainties, affecting frequency channels: 30GHz, 44GHz and 70GHz.

Affects large angular scales. $\ell \leq$ 30, with 44GHz being the major problem.

Planck 2015 data: Large angular scale data, i.e. $\ell \leq 20$ removed. Data from $20 < \ell \leq 40$ high pass filtered. Data did not agree well with FFP8 simulations.

Planck 2018 data: Data from all angular scales present. Systematics issues acknowledged and modelled in FFP10 simulations.

Planck 2015 Results

ℓ_{min}	ℓ_{max}	$S_{H}^{EE} \times 10^{-13} [\text{in K}^2]$	r _{EE}	$\hat{\lambda}$	p-value
40	100	6.94	0.031	$(l = 268^\circ, b = 56^\circ)$	0.37
40	125	6.92	0.029	$(l = 268^\circ, b = 56^\circ)$	0.57
50	100	6.36	0.036	$(l = 260^\circ, b = 57^\circ)$	_
50	125	6.36	0.036	$(l = 260^\circ, b = 57^\circ)$	_



Planck 2015 Results



Planck 2018 Results

Мар	ℓ_{min}	R estimator			D estimator		
		А	Â	p-value	А	Â	p-value
Comm.	10	0.618	$(l = 352^{\circ}b = -13^{\circ})$	< 1/300	0.584	$(l = 350^{\circ}b = -14^{\circ})$	< 1/300
	20	0.295	(l = 346^{\circ}b = -18^{\circ})	0.15	0.326	(l = 347^{\circ}b = -20^{\circ})	0.09
	40	0.286	(l = 352^{\circ}b = -13^{\circ})	< 1/300	0.279	(l = 352^{\circ}b = -13^{\circ})	< 1/300
SMICA	10	0.197	$(l = 228^{\circ}b = -1^{\circ})$	0.37	0.182	$(l = 226^{\circ}b = -2^{\circ})$	0.44
	20	0.078	(l = 232^{\circ}b = -30^{\circ})	0.83	0.076	$(l = 233^{\circ}b = -31^{\circ})$	0.84
	40	0.069	(l = 53^{\circ}b = 75^{\circ})	0.65	0.069	$(l = 58^{\circ}b = 75^{\circ})$	0.65
SEVEM	10	0.587	$(l = 352^{\circ}b = -21^{\circ})$	< 1/300	0.596	$(l = 350^{\circ}b = -22^{\circ})$	< 1/300
	20	0.581	$(l = 346^{\circ}b = -4^{\circ})$	< 1/300	0.576	$(l = 344^{\circ}b = -5^{\circ})$	< 1/300
	40	0.473	$(l = 345^{\circ}b = 0^{\circ})$	< 1/300	0.476	$(l = 344^{\circ}b = -1^{\circ})$	< 1/300
NILC	10	0.374	$(l = 332^{\circ}b = -16^{\circ})$	0.77	0.389	$(l = 333^{\circ}b = -17^{\circ})$	0.75
	20	0.258	$(l = 329^{\circ}b = -29^{\circ})$	0.94	0.231	(l = 332^{\circ}b = -29^{\circ})	0.97
	40	0.351	$(l = 2^{\circ}b = -1^{\circ})$	0.41	0.347	(l = 3^{\circ}b = -2^{\circ})	0.44



Summary

- Planck temperature data has confirmed the presence of the dipole modulation of the CMB temperature signal with an amplitude of 0.07 along $(232^\circ, -14^\circ)$ at about 3σ .
- The dipole modulation is scale dependent and is only present up to $\ell \sim$ 60.
- The exact form of the scale dependence is unknown.
- Due to the large residual systematics in the Planck polarization data, there hasn't been any detection of power asymmetry in Planck polarization data.
- Future of the hemispherical asymmetry study would be to use CMB polarization with CMB temperature to improve understanding of the effect.

Thank you!